The Case for Quantum-Coherent Nuclear Engineering

Florian Metzler, PhD 2/7/2021

Confidential

Overview

- · Quantum 2.0 and Nuclear: Programmatic considerations
- So what is coherence?
- · What can coherence do for us? An overview of quantum dynamics
- The future of quantum dynamics simulations
- · How can weak couplings have large effects?
- · Building intuition on dynamics and transition rate changes
- · Quantum-coherent engineering
- · Quantum coherence at the nuclear level
- · Nuclear reactions and nuclear state transitions
- What's next?
- How to create quantum-coherent nuclear ensembles?
- · Reviewing relevant experiments
- Proposed experiments (toward definitive confirmation and applications)
- · What does quantum-coherent nuclear engineering look like?

Quantum 2.0 and Nuclear: Programmatic considerations

Two ways to approach research

DEMAND SIDE

What problems need solving? What do we know that might help address them?

FORTHCOMING ARTICLE:

"Nuclear fusion rate enhancement in solid-state environments"

SUPPLY SIDE

What new tools are emerging? In what ways can they be applied?

PROPOSED ARTICLE:

"The Case for Quantum-Coherent Nuclear Engineering"

Two ways to approach research

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- Atomic physics
- · Nuclear physics
- Quantum dynamics

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PROPOSED ARTICLE:

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Quantum 2.0



Nuclear science



What is Quantum 2.07

A definition



"Quantum 2.0 refers to the development and use of many-body quantum superposition, entanglement, and measurement to advance science and technology. Examples are quantum computing and simulation, quantum communications, and quantum sensing.

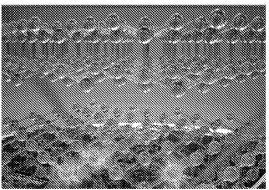
New resulting technologies will potentially go for beyond the (quantum 1.0) capabilities offered by systems without the conceptual need for large-scale superposition or entanglement, examples of which are conventional semiconductor electronics, loser-based communication systems and magnetic-resonance medical imagers."



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More on Quantum 2.0: Basic Energy Sciences Roundtable

Opportunities for Basic Research for **Next-Generation Quantum Systems**



Participants:

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More on Quantum 2.0: Basic Energy Sciences Roundtable

Emphasizing controlled superposition

"From the transistor to the molecular switch, quantum mechanics is at the heart of nearly all materials properties and chemical processes. Still, the subtleties of quantum behavior are often hidden from view. As a result, only a small number of scientific techniques and technological applications take advantage of the unique phenomena of quantum superposition and entanglement. Harnessing these counterintuitive properties of matter promises to yield revolutionary new approaches to computing, sensing, communication, and metrology. [..]

Quantum-coherent systems have been discovered that exhibit remarkable properties and ever increasing coherence times. However, understanding of how these systems interact (for example, the complexity of entanglement) is currently limited to a small number of systems. Advances in tahis field require an understanding of the scaling of coherence lengths and times with system size and complexity, and the identification of new signatures of quantum states in artificial quantum-coherent systems."

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More on Quantum 2.0: Basic Energy Sciences Roundtable

Coherent quantum-to-quantum transduction

"[Goals:] Discover novel approaches for quantum-to-quantum transduction: The coherent transduction of information [and energy!] from one modality to another, at the single-particle or quantum level, is at the core of quantum measurement and information processing. The development of quantum science will also contribute to areas outside of computing. Control over quantum-coherent states in artificial systems may lead to enhanced transduction for novel electronics, efficient light harvesting and photovoltaics, new techniques for cosmology and nuclear science, and sensing capabilities that are orders of magnitude beyond current standards."

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More on Quantum 2.0: Basic Energy Sciences Roundtable

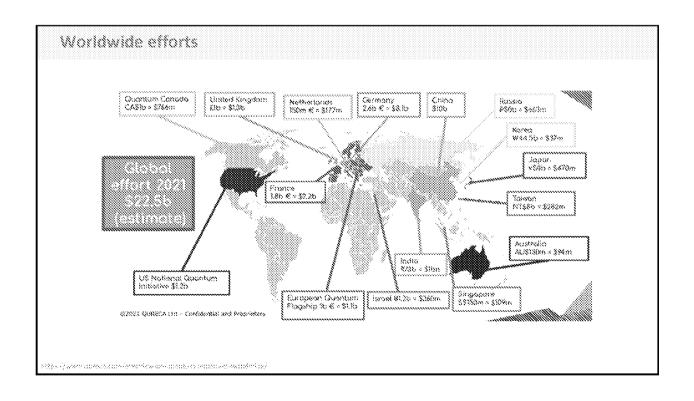
Photosynthesis, strong coupling, tunable lattices

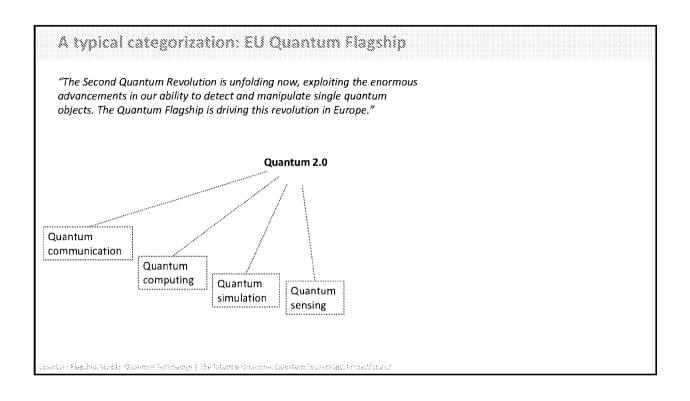
"The light harvesting apparatus of plants and bacteria possesses remarkable capabilities of reaching near unit efficiency for the conversion of an absorbed photon to an electron which goes on to initiate the [..] chemical reactions of photosynthesis. [..] Recent ultrafast spectroscopy and theoretical work have revealed insight into the quantum dynamics of the energy transfer and show how the excitonic degrees of freedom interact cooperatively with phonons to ensure the optimality and robustness of the transport in the presence of disorder. This approach can provide important insights for the design of quantum-to-quantum transducers. [..]

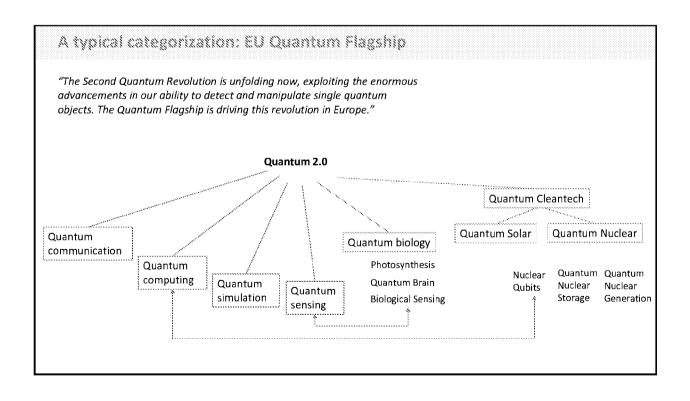
Theory of quantum state transfer: A key challenge in any quantum transduction scheme is to find a quantum system that is flexible enough to couple strongly and coherently to a variety of distinct systems and is thus capable of acting as a "quantum bus." [..] A prime example is the use of vibrational modes of a mechanical resonator as a quantum bus. [..]

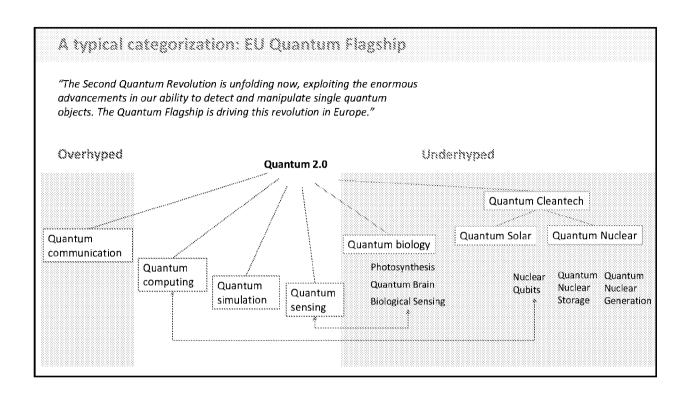
Arrays of dopant atoms, or metal-organic sites, with specified arrangements can be introduced to create highly tunable superlattices to enhance particular phonon, plasmon, or optical modes."

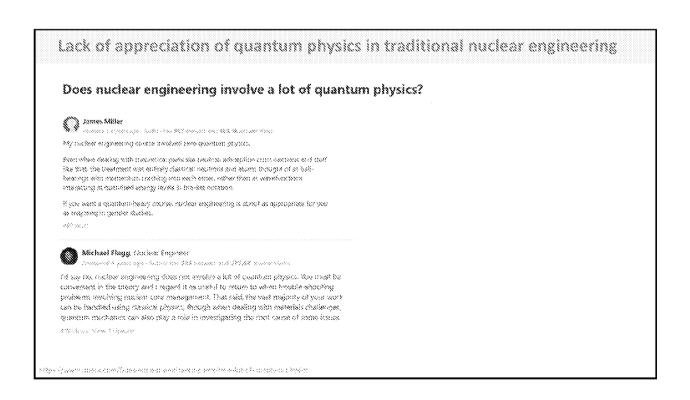
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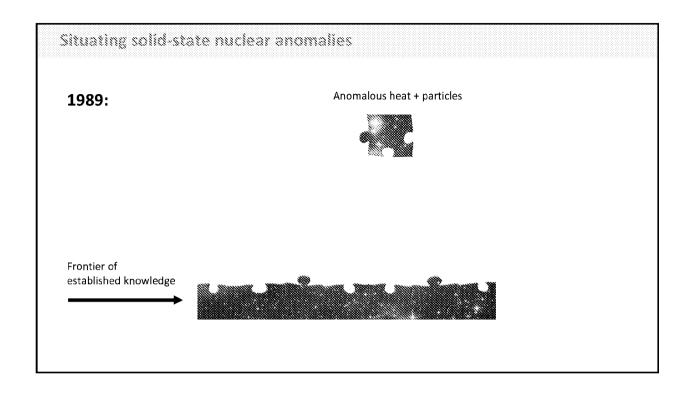


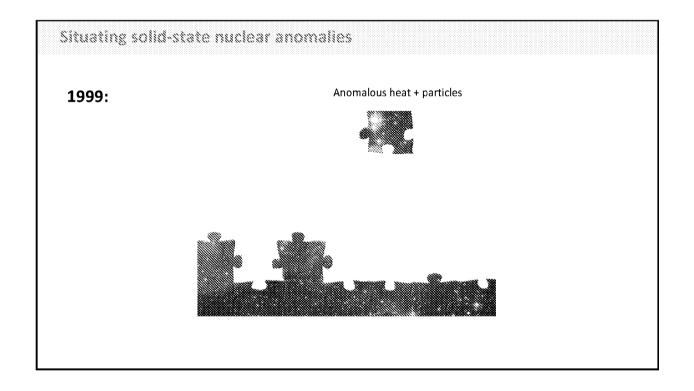


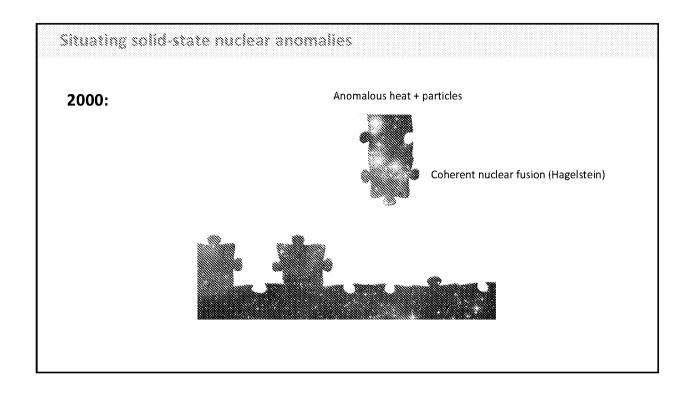


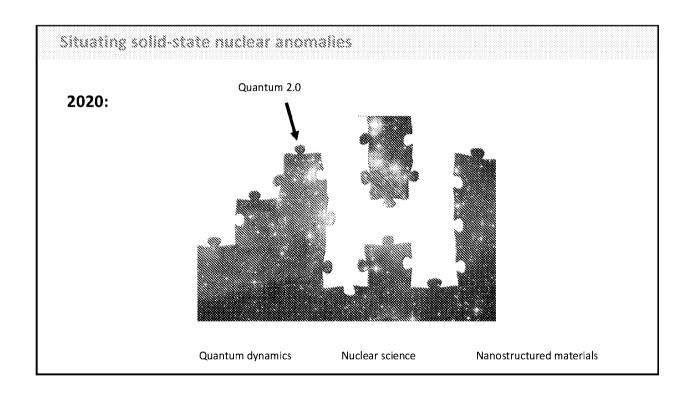


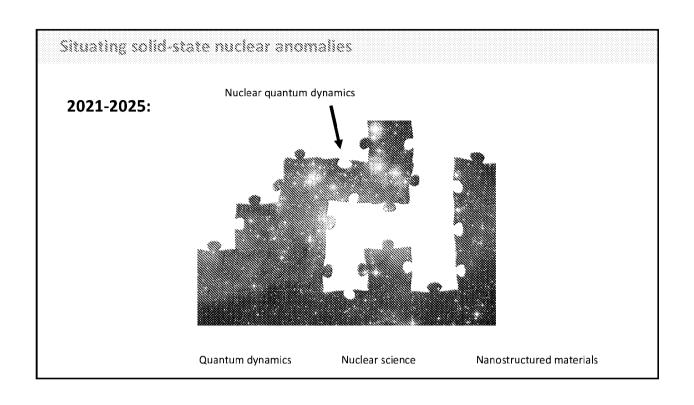


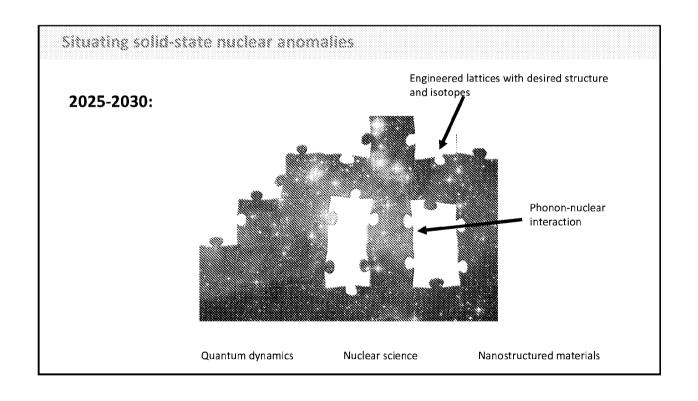


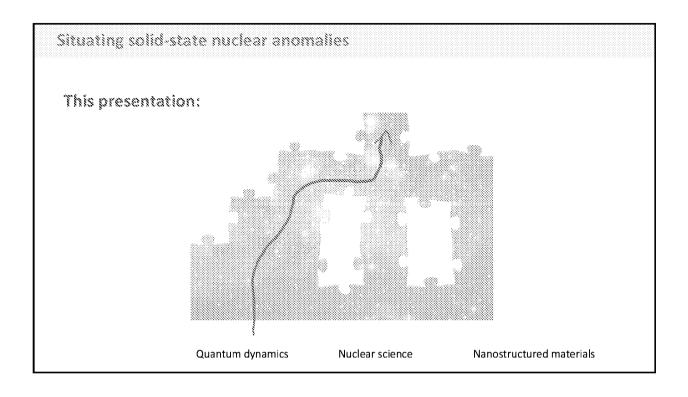












A timely frontier

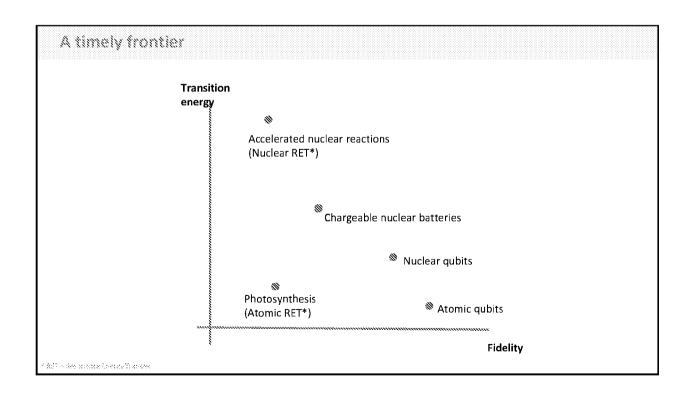
Outlook

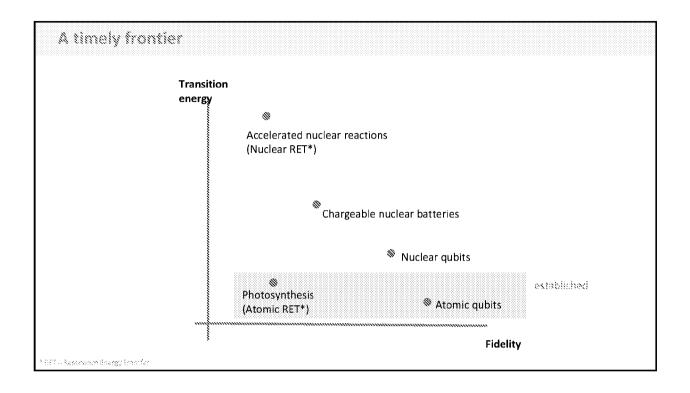
I will show the gradual transition from quantum computing to quantum-coherent nuclear engineering.

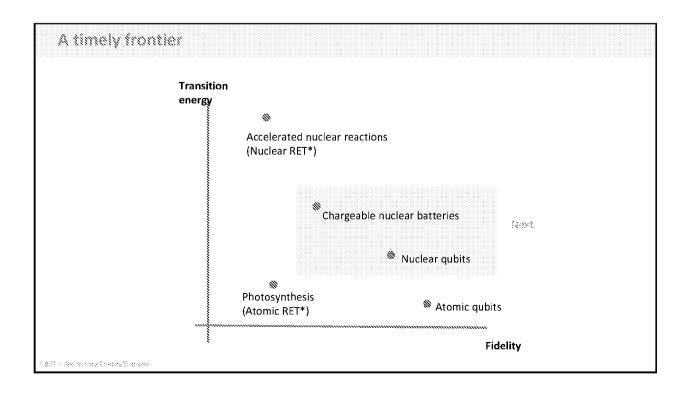
Spoiler:

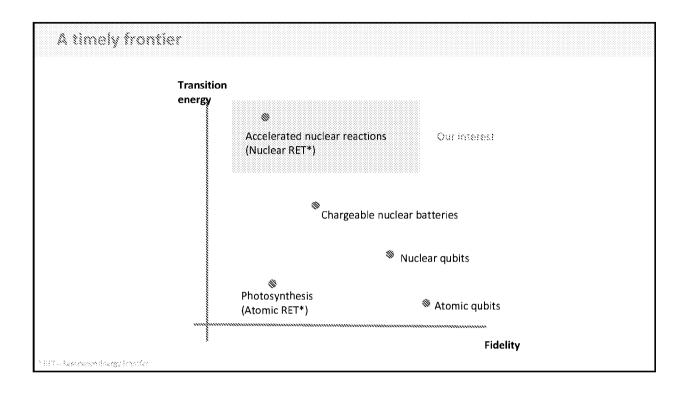
Quantum-coherent nuclear engineering is essentially quantum computing with high-energy qubits!

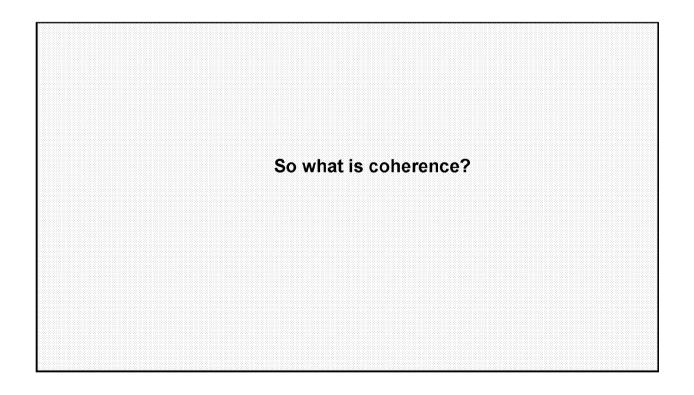
A timely frontier The fidelity/energy tradeoff Compare with classical electromagnetics: • A low-energy electromagnetic wave constitutes a signal • A high-energy electromagnetic wave constitutes energy transfer











A definition:

A quantum system sufficiently stable so that it can maintain discrete energy/information states, including superposition with coupled subsystems (i.e. energy/information is held collectively by multiple subsystems).

→ allows for precise control of individual quantum states

The opposite is decoherence:

Irreversible dissipation of energy/information into the environment, akin to heat losses in thermodynamics.

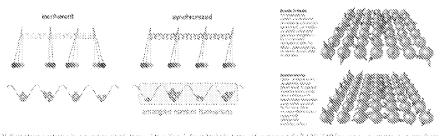
→ allows only to work with statistical aggregates (Quantum 1.0)

Introduction to coherence

More concretely:

To achieve this stability, need:

- Strong coupling internally (between discrete states)
- · Weak coupling to the environment (to infinite states/continuum)
- Phase synchronization (often follows from strong coupling)

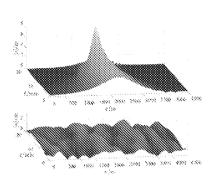


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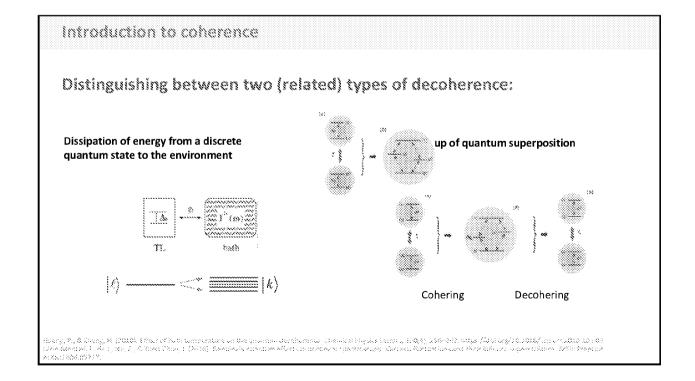
Another way of putting it:

There is a degree of order in the system that allows for the occupation of a few high-energy states vs. many lower-energy states.

Compare with classically coherent wave mechanics:



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We can think about the degree of coherence on a spectrum

Energy is completely distributed across environment

Energy is stably held in a single quantum system

Energy is held collectively through a superposition of quantum systems

Characteristics:

strong coupling to outside

weak coupling to outside

weak coupling to outside strong coupling within

Examples:

Quantum 1.0

Nuclear decay

Qubits

BEC, photosynthesis

Incoherent nuclear engineering



Coherent nuclear engineering

Introduction to coherence

From an exotic state towards a common tool

Previous belief:

Coherence is an exotic and extremely rare state, only possible under exceptional circumstances (close to OK, very short lived).

Increasingly common realization:

Coherence shows up in all kinds of places, especially in nature. It often can be stabilized by boundary conditions and can be surprisingly longlived.

Example: growth of coherence lifetimes in qubits

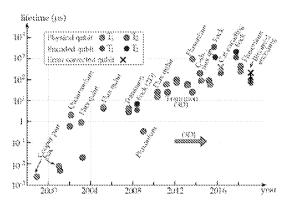
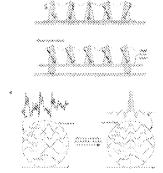


Figure 3.6: Evolution of lifetimes and experence times in superconducting publis.

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Inducing coherence through couplings

Formation of coherence as synchronization





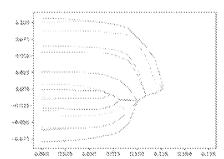
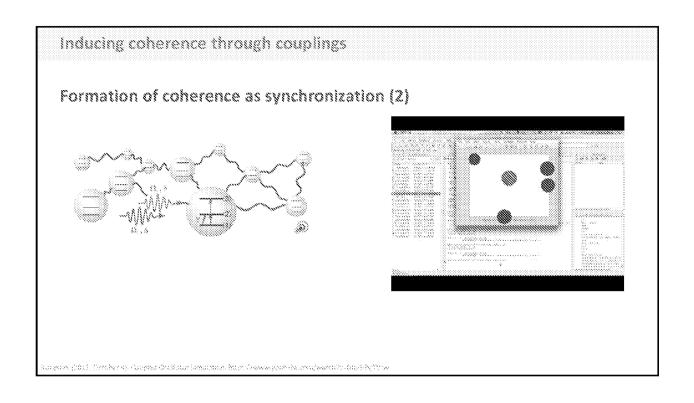
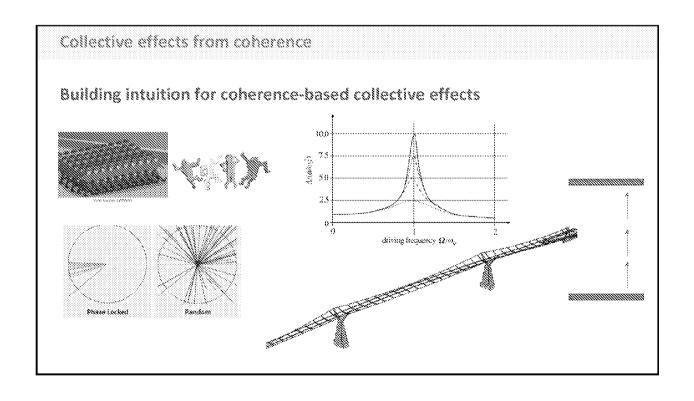


Fig) 9 The Kuramoto model for 75 Poincare oscillators showing the frequencies as a function of the coupling coefficient)

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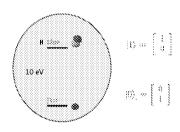


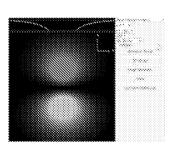


What can coherence do for us? An overview of quantum dynamics

Quantum dynamics overview

Excited states in the absence of an environment are stable oscillators

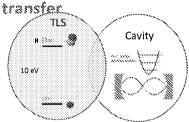


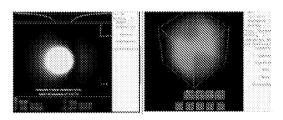


$$\begin{split} \hat{H}_M &= \begin{bmatrix} \frac{\partial r_0}{\partial x} - \frac{\partial}{\partial x_0} \\ \\ \frac{\partial r_0}{\partial t} \psi &= \hat{H} \varphi \\ \\ \psi(\mathbf{r}, t) &= e^{-iE_{c}t/\hbar} \, \psi(\mathbf{r}) \end{split}$$

$$P(\mathbf{r},t) = |\psi(\mathbf{r},t)|^2 = |\psi(\mathbf{r})|^2$$

A coupling can lead to discrete energy





 $\sigma_{\rm h}^{\rm L}(0)/(1)$. As an decays time (1) \rightarrow (0) and only a shorter (in the $\sigma^{\rm th}$ mode).

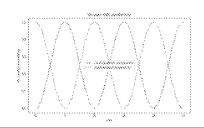
out [1] (2) Attends a mixed frice [1] of [1] and absents a photoe (from the order).

 $a_{\rm M}^{\rm H}(0)$. After it substitute (time (0) \sim [1) and emits a photon (in the $m^{\rm th}$ mode), $a_{\rm M}^{\rm H}(0)$. Along decays from (1) \sim [8) and about a photon (from the $m^{\rm th}$ mode),

$$\mathcal{H} = \mathcal{H}_0 + V = \frac{\hbar}{2}\omega_0\sigma_x + \sum_m \hbar\omega_m \left(\sigma_m^i\sigma_m + \frac{1}{2}\right) + V$$

LS Electromagnetic field Coupling strength

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Quantum dynamics overview

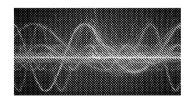
The option for many discrete energy transfer pathways leads to exp. decay

$$\dot{H} = \hat{H}_M + \dot{H}_R + \dot{H}_I$$

Radiation field as an infinite set of oscillators:

$$\hat{H}_{\mathcal{B}} \ = \ \sum_{\mathbf{k},\sigma} \, \hbar \omega_{\mathbf{k}} \, \left[\, \hat{a}^{\dagger}_{\mathbf{k},\sigma} \hat{a}_{\mathbf{k},\sigma} \, + \, \frac{1}{2} \, \right]$$

With coupling: $\hat{H}_{I} = -\hat{\mathbf{d}} \cdot \hat{\mathbf{E}}(\mathbf{r})$



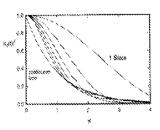
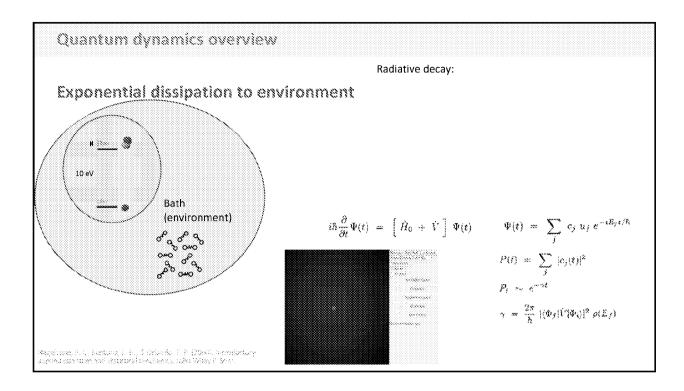
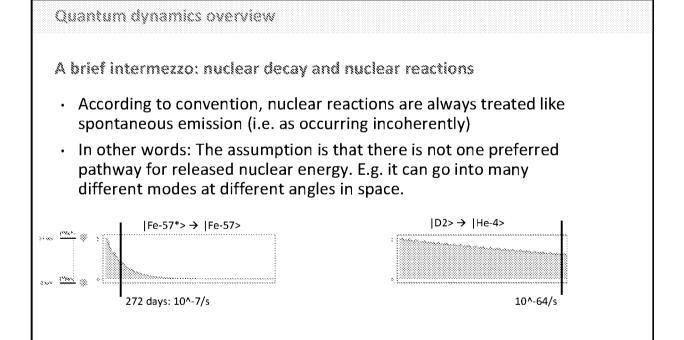


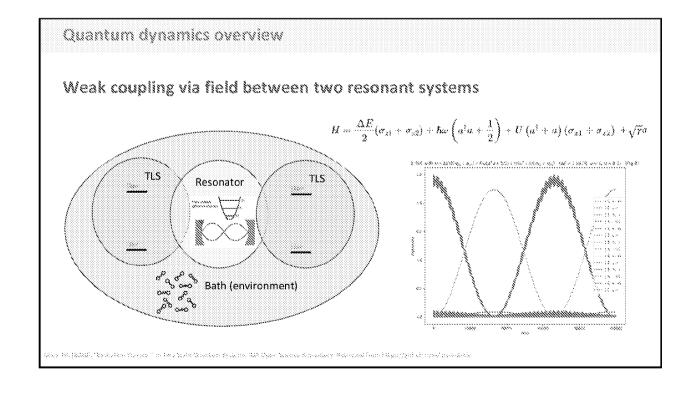
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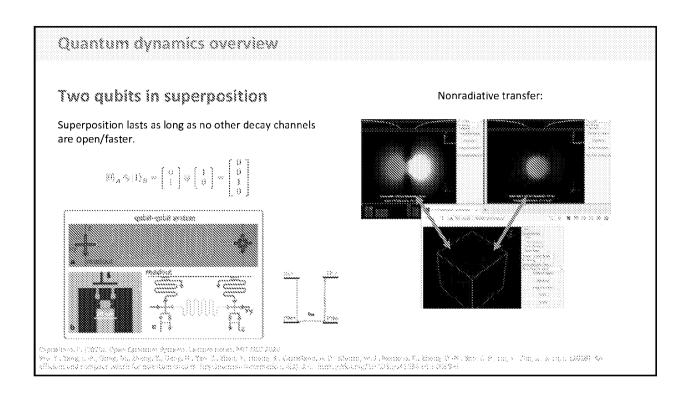
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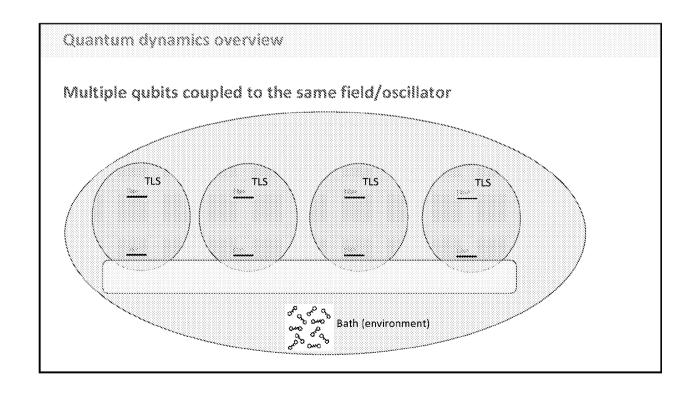




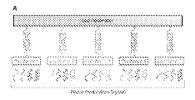
Coupling to field with dissipation $H_{Raii} = \omega_r(a^*a + \frac{1}{2}) + \frac{\alpha_s}{2}\sigma_z + g(a^* + a)(\sigma_+ + \sigma_-) + \sqrt{r}a$







Multiple qubits coupled via a quantum bus



$$H=H_0+V, \quad V=\frac{\mathcal{U}}{2}\sum_{i=1}^N d_i^0 d_i \Big(d_i^0 d_i+1\Big).$$

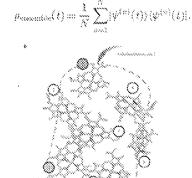
If you have multiple qubit gates, can have superposition of many systems (less relevant for quantum computing but for other techniques such as quantum annealing).

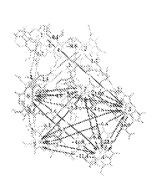
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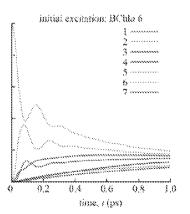
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Quantum dynamics overview

Formation of an ensemble: superposition of multiple subsystems







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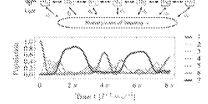
Photosynthesis

$$\begin{split} H_{tot} &\approx H_{tops} + H_{core} + H_{tot}, \\ H_{tops} &\approx \frac{h_{tot}}{2} \delta_2 + h_{to}, \\ H_{core} &\approx \sum_{\mathbf{k}} h_{tot} \delta_{\mathbf{k}}^2 \delta_{\mathbf{k}}, \\ H_{tot} &\approx \sum_{\mathbf{k}} \delta_2 \otimes \left(h_{tot} \delta_{\mathbf{k}}^2 + h_{tot} \delta_{\mathbf{k}}^2 \right). \end{split}$$

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Quantum dynamics overview

Phonon-mediated exciton transfer



$$\hat{H} = \hat{H}_{c} + \hat{H}_{ph} + \hat{H}_{c+ph}.$$

$$\hat{H} = \sum_{i=1}^{N} \delta\left(\Omega_{i} + f_{i}^{2} \omega^{2}\right) \left\{i\right\} \left\{i\right\} + \sum_{i=1}^{N} \delta I_{i} \left\{i\right\} \left\{i\right\}$$

$$\hat{H}_{i,n} = \hbar \omega \left\{i, \hat{h} + \frac{1}{2}\right\}$$

$$\hat{H}_{i,n} = \sum_{i=1}^{N} \omega_{i} \left\{i\right\} \left\{i\right\} \approx \int_{0}^{1} \omega_{i} \left\{i\right\} \left\{i\right\} \left\{i\right\} \left\{i\right\}$$

VIDEO: accelerated state transitions due to genetically engineered spacing of excitons:

https://www.youtube.com/watch?v=91vhoxR1Lts



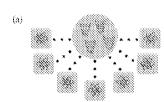
$$H_S = \sum_{n=0}^{N} \epsilon_{s,n} a_{s,n}^{\dagger} a_{s,n} + \sum_{n=\infty}^{N} V_{s,n} (a_{s,n}^{\dagger} a_{s,n} + a_{s,n}^{\dagger} a_{s,n})$$

 $V_{mn} \sim \frac{1}{2} \left(p_{mn} \cdot p_{mn} \cdot \frac{1}{2} \left(p_{mn} \cdot \mathbf{B}_{mnn} \right) \left(p_{mn} \cdot \mathbf{B}_{mnn} \right) \right)$

Multiple qubits interacting via multiple field modes: Spin-boson model

More specifically we consider a system of N qubits interacting with K bosonic modes (Fig. 1(a)) by means of a spin-boson Hamiltonian Massimo Paina at al. (2006) We at al. (2003) of the form

$$H = \sum_{i=1}^{N} \frac{\Omega_i}{2} \sigma_s^i + \sum_{k=1}^{K} \omega_k \theta_k^{\dagger} b_k + \sum_{i,k} \lambda_{i,k} \sigma_s^i (b_k^{\dagger} + b_k), \omega$$



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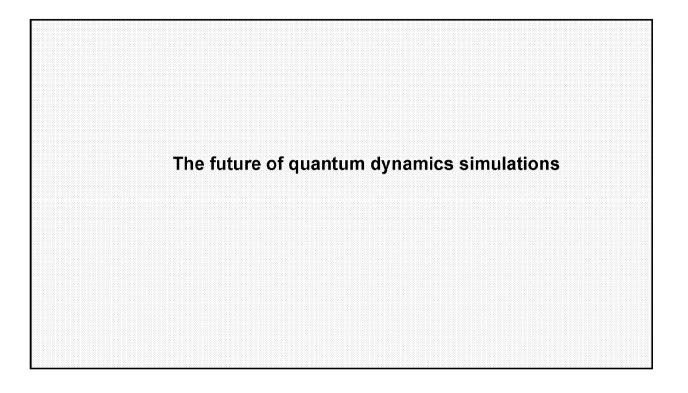
Modelling the full picture

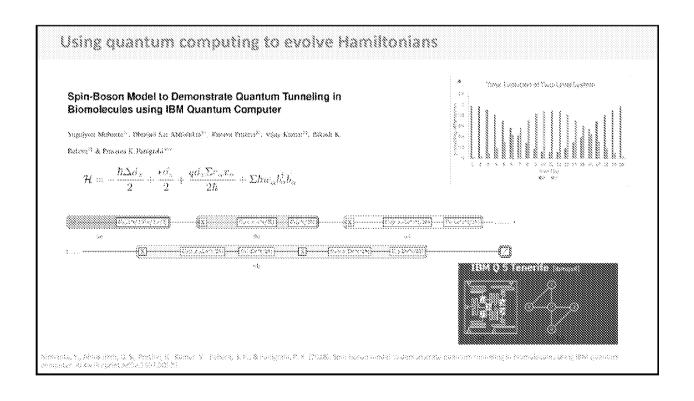
Sneak preview

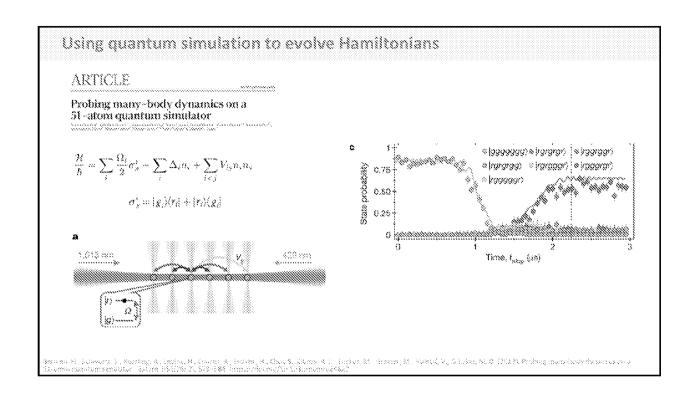
This is the model for quantum-coherent nuclear engineering:

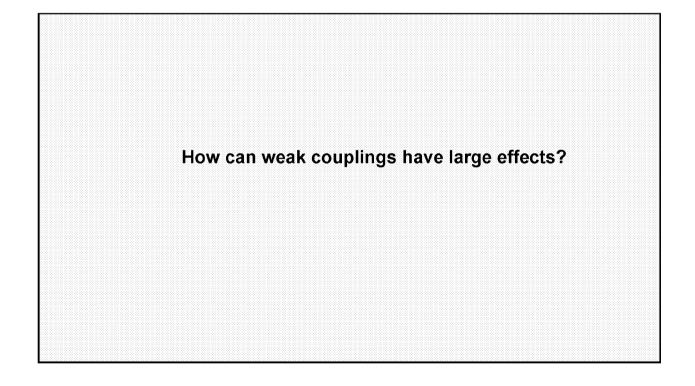
$$\begin{split} \hat{H} &= \sum_{i} \mathbf{M}_{i} c^{2} + \sum_{\mathbf{k},\sigma} \hbar \omega_{\mathbf{k},\sigma} \hat{a}_{\mathbf{k},\sigma}^{\dagger} \hat{a}_{\mathbf{k},\sigma} + \sum_{i} \sum_{\mathbf{k},\sigma} \mathbf{V}_{j,\mathbf{k},\sigma} \\ \\ \mathbf{V}_{j,\mathbf{k},\sigma} &= \mathbf{a}_{i} + \left(\frac{\partial \mathbf{P}_{i}}{\partial a_{\mathbf{k},\sigma}^{\dagger}} \hat{a}_{\mathbf{k},\sigma} + \frac{\partial \mathbf{P}_{i}}{\partial a_{\mathbf{k},\sigma}^{\dagger}} \hat{a}_{\mathbf{k},\sigma}^{\dagger} \right) \end{split}$$

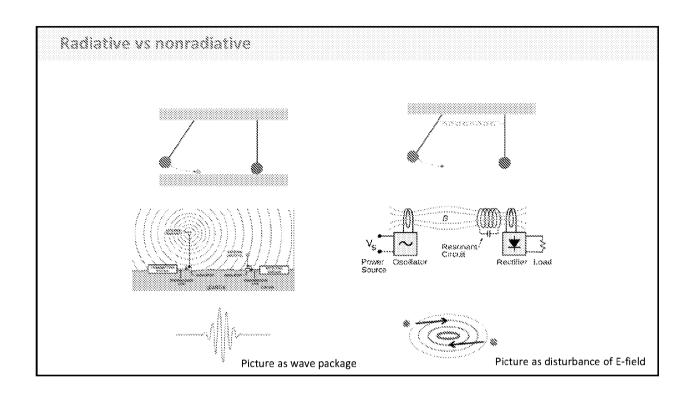
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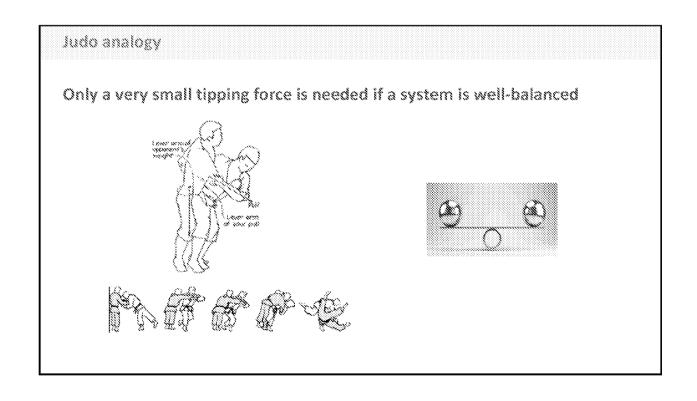


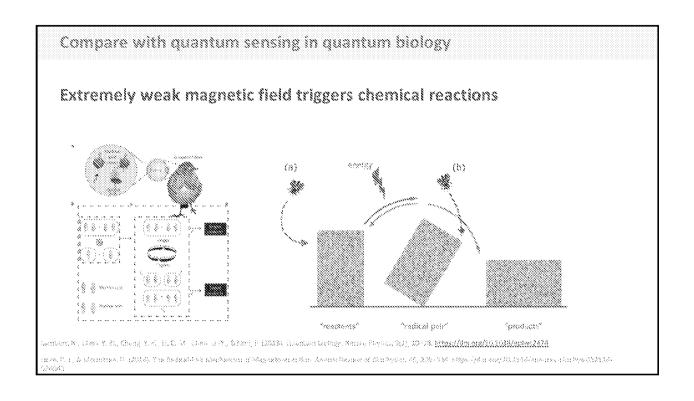


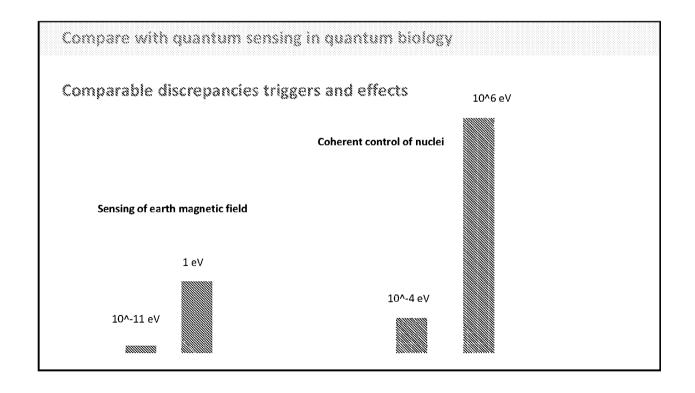












Building intuition on dynamics and transition rate changes

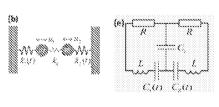
Dynamics and transition rate changes

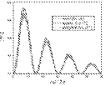
Classical dynamics can reproduce quantum dynamics precisely

PSSYSTEM, SUNDEW ARS, (SCIENCE)

Coherent quantum states from classical oscillator amplitudes

In the first days of construct precisions: Dies promoted as endings between the latest dependence on fillelooks of on exponence of the factorinages required and the chooses of promotes of the factorinages required and the chooses of promotes on the construction of expositions of the factorinages of the factorinages of the factorinages of the construction of the factorinage of the factorinages of the fac





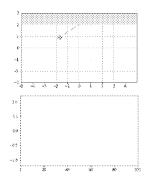
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Dynamics and transition rate changes

Dissipative exponential decay





$$\frac{dz}{dt} = \frac{p(t)}{m}, \quad \frac{dp}{dt} = -m\omega^2 z \quad \cdots; \quad \frac{do}{dt} = -i\omega\alpha(t)$$

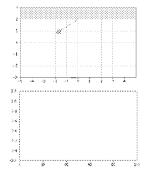
$$\begin{split} \langle X \rangle &= \frac{1}{\sqrt{3}} (\alpha_0 e^{-i\omega t} + \alpha_0^2 e^{i\omega t}) \\ \langle P \rangle &= \frac{1}{\sqrt{3}} (\alpha_0 e^{-i\omega t} + \alpha_0^2 e^{i\omega t}) \end{split}$$

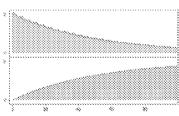
$$X=(a+a^{4})/\sqrt{2}, \qquad P=-i(a-a^{4})/\sqrt{2}, \qquad \mathcal{H}=\omega(a^{\dagger}a+\frac{1}{2}),$$

Hote tray, P. A., & Proton, E.S. G. (1988). Coupled-persolate invaried are translated internative former effect, 2098-8, \$195, 1818–1818.

Dynamics and transition rate changes

Dissipative exponential decay





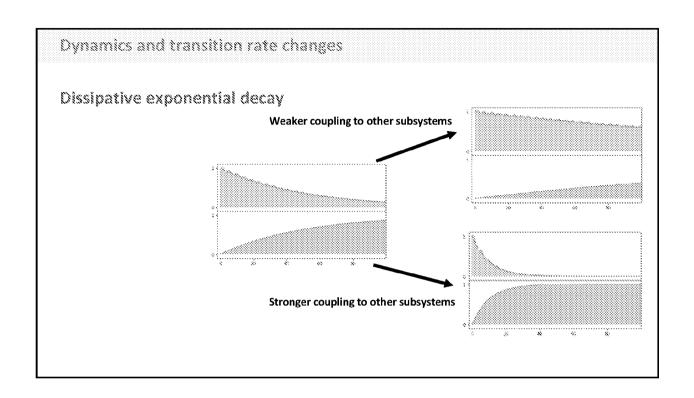
$$i\hbar\frac{\partial}{\partial t}\Psi(t) \;=\; \left[\; \dot{H}_0 \;+\; \dot{V}\; \right]\; \Psi(t)$$

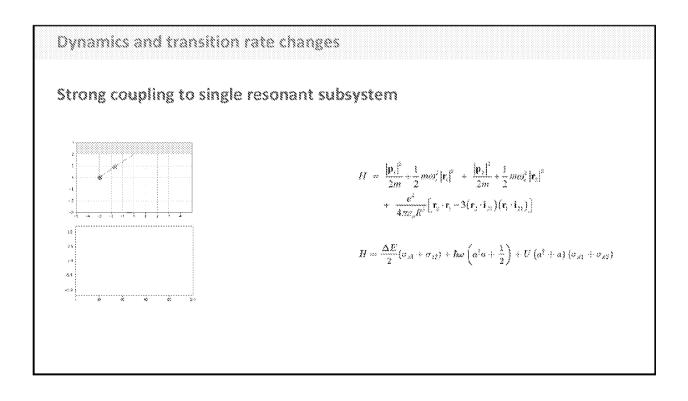
$$\Psi(t) = \sum_{j} c_{j} u_{j} e^{-i\partial_{j}t/\hbar}$$

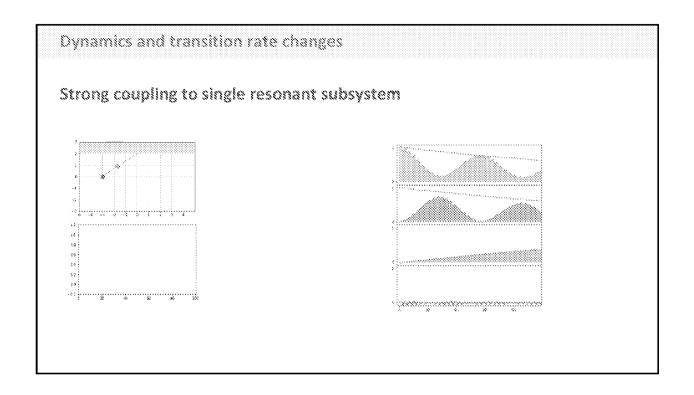
$$|P(t)| = |\sum_{j} |e_{j}(t)|^{2}$$

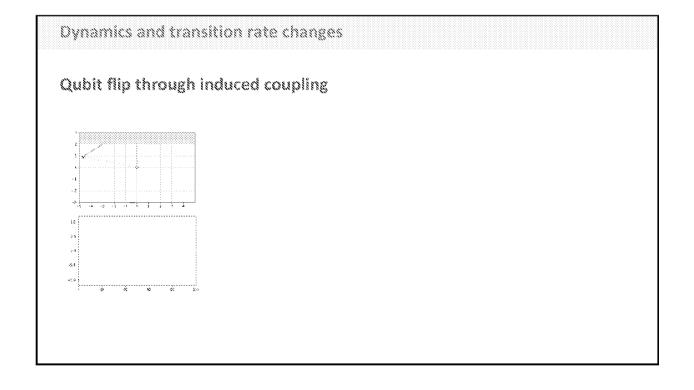
$$P_i \sim e^{-\gamma t}$$

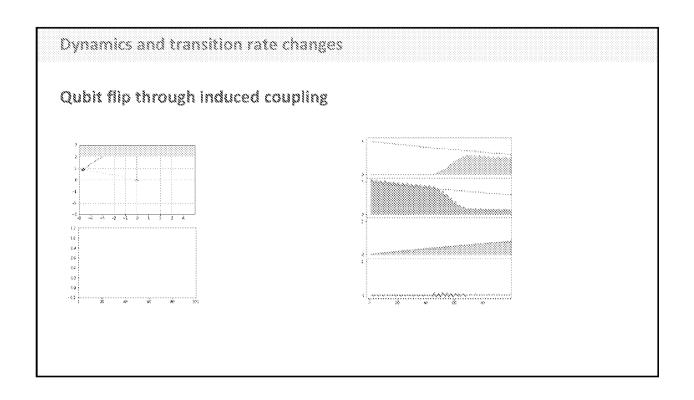
$$\gamma \; = \; \frac{2\pi}{\hbar} \; |\langle \Phi_f | \dot{V} | \Phi_i \rangle|^2 \; \rho(E_f)$$

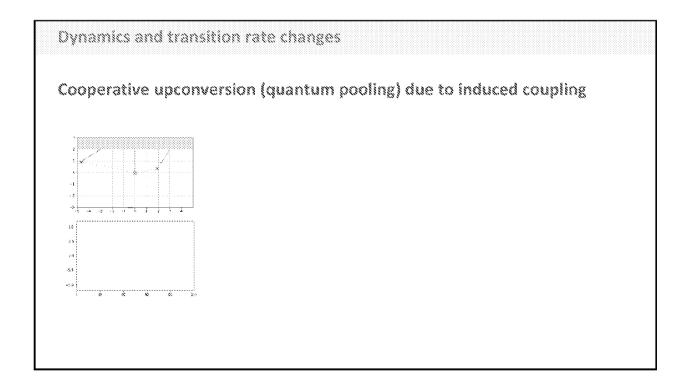






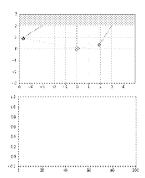


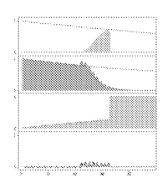




Dynamics and transition rate changes

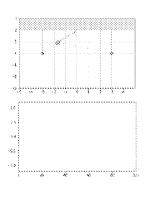
Cooperative upconversion (quantum pooling) due to induced coupling

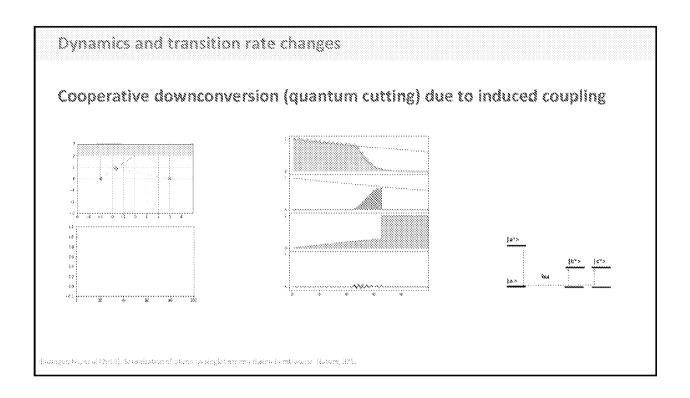


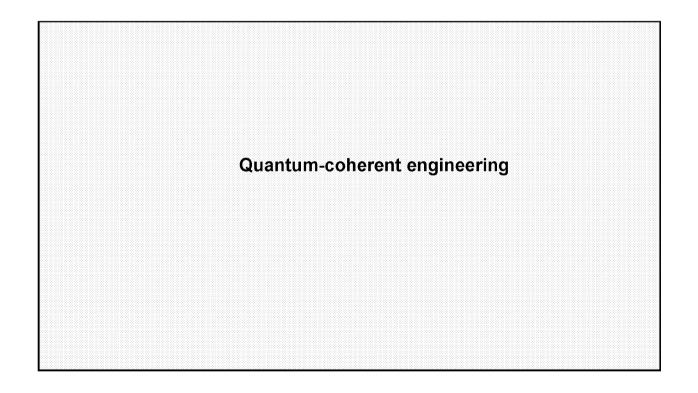


Dynamics and transition rate changes

Cooperative downconversion (quantum cutting) due to induced coupling



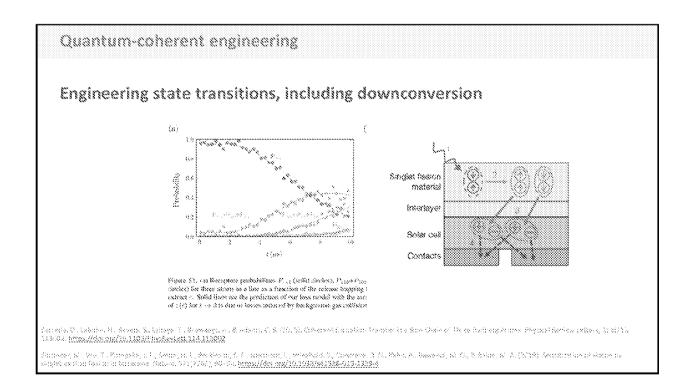


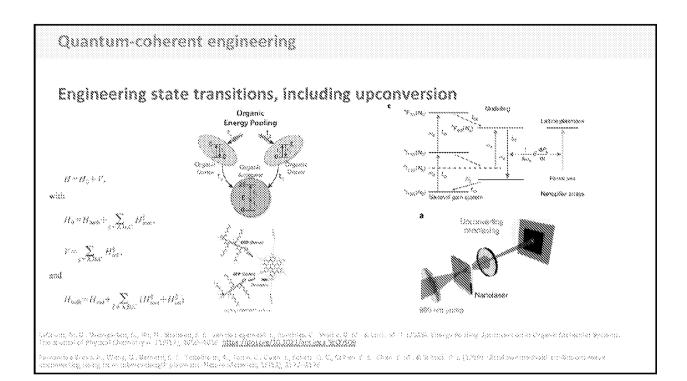


Engineering transition rates Figure 4:14% belongly decay transition of the properties of the properti

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reasson increases)





Quantum-coherent engineering

Predicting dynamic outcomes from system parameters

TABLE 1. Table of transitions between lauthonide dopost ions is crystal field environs: i-energy pooling, ii-two plants APTE, iii-sossistation, iv-quantum corring.

A (Transition)	8 Chanaineo	© (Transilion)	Туре	Kelerovees 18, 23, 28	
685 (P50 2850)	$Yb^{34}(^2F_{gg}\rightarrow^2F_{gg})$	18 ⁵ (² B ₈ ~ ³ B ₈)	;		
$B^{3n}(^2S_{2m},a^2P_{2m})=$	$Y_0^{11} Y_1^2 F_{22} = {}^3 F_{22}$	$4m^{2\phi}CB_{\phi}m^{\phi}G_{\phi})$	8, 8	28, 22	
$O^{3/(2F_{2Q} - e^2F_{2Q})} =$	$Y b^{i+} ({}^{i}F_{0i} \cdots F_{2i})$	$\operatorname{Ex}^{S^{*}}(D_{\theta^{-n}}^{-1}\mathcal{F}_{\phi})$	3	25	
$arphi^* \psi p_j { ightarrow} g_0 = -$	$\mathbf{e}e^{i\phi}(\partial \Omega_{j}\omega^{j}R_{i})$	$8e^{5s}(^3D_2\cdots ^3S_8)$	į.	29, 30	
$\partial V^{*}(F_{SS}, \sigma^{2}P_{SS})=$	YS 32 F 50 - 2 F 50	$\mathcal{W}_{0}^{(3)}(^{5}f_{8}^{5}\mathcal{S}_{3}^{5}\mathcal{F}_{6})$	22	i3, 2i	
$\phi^{\mu_1}(\Phi_{\mu_2} + \Psi_{\mu_2}) =$	$Yb^{10}({}^{1}F_{20},,F_{20})$	$\operatorname{Ex}^{4\phi}({}^{\phi}I_{10,27}^{4}F_{22})$	#	48, 19, 21, 24, 26	
$ abla^{3}(^{3}P_{am}{}^{3}O_{a}) = -$	$\nabla b^{34} (^2 F_{\gamma \gamma \gamma \gamma \delta}^2 F_{\gamma \gamma \delta})$	$B_{\epsilon}^{NN}(^{3}R_{A},,^{3}G_{\epsilon})$:::	36	
$\mathscr{L}^1(Y_{130}\circ Y_{230})=$	$\mathcal{E}m^{3,\alpha}({}^{3}H_{\alpha} + {}^{1}H_{\alpha})$	$Es^{2} + A_{10,0} - A_{No}^{2}$	111	33	
$\{\phi^{g_{g_{1}}}(^{1}P_{g_{2}},,^{g_{g_{g_{g}}}}\}_{i,j}\}$	$Yb^{3,0}({}^{2}F_{2,0}{}^{3}F_{2,0})$	$85^{35}({}^{3}I_{8}{}^{-1}{}^{5}S_{5},{}^{5}F_{3})$	308	35	
$\langle \phi^{1} \rangle \langle {}^{3}F_{a} \cdots {}^{3}H_{a} \rangle$	$\nabla b^{3N} e^{N} F_{m_2 \cdots m_N} F_{m_N}$	Tso Cas W.	;; ;	32, 38	
$\operatorname{Vd}^{2,r}({}^{\alpha}{\mathcal F}_{22}+{}^{\alpha}{\mathcal F}_{122})=$	$Y_8^{j+}({}^3F_{jj} - {}^2F_{j,j})$	$183\%(^{3}O_{2} \rightarrow ^{2}O_{1})$	iii	39	
$\mathfrak{h}^{38}({}^{5}F,^{5}D_{8})$	$\operatorname{Ext}^{S_1}({}^2F_1,,{}^5H_j)$	$\Theta S^{3,0}({}^{\circ}O_{S^{-1}}{}^{3}S_{S^{2}})$	Ŝŵ	43	
$\Psi^{\gamma}({}^{8}D_{\gamma} + {}^{2}F_{\gamma})$	$Ee^{3/4}(^4S_{Np}, e^4J_{18/2})$	Ex 14 (4/1653 - 43 50)	įsv.	3.2	
ويعود سايهوا والمعاور	860 (385 m 85)	$2N^{2/3}({}^3S_8 + {}^3S_3)$	èv	43	

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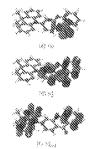
Quantum-coherent engineering

Predicting dynamic outcomes from system parameters

Table 1. Vectical Absorption ΔE_{∞} and Emission ΔE_{∞} Emergies (eV) as well as Adiabatic Excitation Emergies ΔE_{∞} , (eV) of $\Delta CRKTO^{\infty}$

steis.	હ	stoone somiare	1977.0363.8	35 9100.1	füle f 835.79	28%		 Pabas	
1.5		(0.933.03)	9.00	8383	6.00			2.88	
CA	C)	(B. 1876) 44 - 45 (B. 1876)	3,09	3.41	3.60	2.59	3,27	22,933	4 8, 10 5
\$5.	33.8	(6.36) og ov 10)	363	3.03	3.60%	5.23	2.59	3.85	~ 1/3***
255	1.83	(0.20) 44 47	3.35	4.31	3.82	3,69	3.55	1,60	023442
6.85	129	$(0.269 \phi_{\rm K} \cdot \phi_{\rm c} \phi_{\rm c})$	300	3.54	3.83	3.58	3.71	2174	
VA'	3.23	(0.88) z ₀₀₀ x((0.800 s ₀₀₀ x(3.28	722	3.53	3.5%	3.87	0,28	
glass	15.6	(0.38) x2 x2.4 (0.38) x3 x2.4	3.58	523	3538			2.59	
\$15"	33.5	(0.50) x ₀ ↔ x0	3.52	3.65	837	5.888	261	5,385	

The electric thole moments a Orders) and modeltor densities for over competed at the generalizes individual $(S_1, siste(S_1, S), S \in V)$ is a CT state to TODET(BSSNP) calculation, front is S_1 at 495 of (DETARCOS) and at 426 of in SECCL $(S_2, siste(S_1, S), S \in V)$ is a CT state in TODET(BSSNP) calculation, not found among the lesson are exceed singles rates in DETARCOS S_2 and S_3 (DSS of) is a CT state in TODET(BSSNP) calculation, not found among the lesson are exceed singles rates in DETARCOS S_3 and S_4 (CC).



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Quantum-coherent engineering

Engineering structures with optimized parameters for desired dynamics and outcomes ("energy transfer editing")

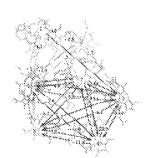
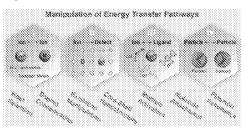


Figure 2



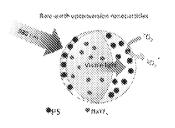


Figure 2. Schemelia illustration entowing four major energy transfer pathways dominated in upconverting nacosystems. A list of chemical and physical approaches to energy transfer manipulation is included in the advance.

a, X. D. L. W., V. K. B. K. (2018). Perty-tunks administrative members that carrylase a Welton memory against as 40 Central Sister, \$12, 25–26.

Summary			
 So it's all about deliberately moving energy/information around in 			
coupled systems.			
Quantum coherence at the nuclear level			

Quantum-coherent engineering

Accelerating state transitions via couplings: nuclei Fig. 3 | Monounterest of Robinsolitations. Personal responses of the sample, showing clearly visible shabit collitations. The libertest at carbon are a faster translation of the course of the sample, showing clearly visible shabit collitations. The libertest at carbon are a faster translation of the course of the about collitations. The libertest at carbon are a faster translation of the storage consulted refere friendly desired crasses of the storage consulted and an expectationally desired crasses of the storage consulted course of the storage consulted and an expectationally desired crasses of the storage consulted course of the storage consulted and an expectationally desired crasses of the storage consulted course of the storage course of the storage consulted course of the storage course of

